

# MATH 514 Comprehensive Exam

Q.1 (10 points) Evaluate the integral  $\int_{-\infty}^{\infty} \frac{\cos x}{(x^2+1)^2(x^2+4)} dx$  using residue theory.

Hint: Consider  $f(z) = \frac{e^{iz}}{(z^2+1)^2(z^2+4)}$ .

Sol: Consider  $f(z) = \frac{e^{iz}}{(z^2+1)^2(z^2+4)}$ ,  $z=i$ ,  $z=2i$   
in the upper half plane.

$$\text{Res}[f(z); i] = \frac{1}{1!} \lim_{z \rightarrow i} \frac{d}{dz} \left[ \frac{e^{iz}}{(z+i)^2(z^2+4)} \right]$$
$$= \lim_{z \rightarrow i} \frac{ie^{iz}(z+i)^2(z^2+4) - e^{iz}(2(z+i)(z^2+4) + 2z(z+i)^2)}{(z+i)^4(z^2+4)^2}$$

$$= \frac{i e^{-1}(-4)(3) - e^{-1}(2(2i)(3) + 2i(-4))}{(16)(9)}$$

$$= \frac{-12 e^{-1}i - 12i e^{-1} + 8 e^{-1}i}{(16)(9)} = \frac{-16 e^{-1}i}{(16)(9)}$$

$$= \frac{-e^{-1}i}{9}$$

$$\text{Res}[f(z); 2i] = \lim_{z \rightarrow 2i} \frac{e^{iz}}{(z^2+1)^2(z+2i)} = \frac{e^{-2}}{9(4i)}$$

$$= \frac{-e^{-2}i}{36}$$

$$\int_{-\infty}^{\infty} \frac{\cos x}{(x^2+1)^2(x^2+4)} dx = 2\pi i \left[ \frac{-e^{-1}i}{9} - \frac{e^{-2}i}{36} \right]$$

$$= 2\pi \left[ \frac{1}{9e} + \frac{1}{36e^2} \right]$$

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Q.2 (10 points) Use Laplace transform to evaluate the integral

$$\int_0^{\infty} \frac{\sin(xt)}{x(x^2+4)} dx$$

Sol:

$$\text{Let } f(t) = \int_0^{\infty} \frac{\sin(xt)}{x(x^2+4)} dx$$

Taking Laplace transform, we get

$$F(s) = \int_0^{\infty} \frac{1}{x(x^2+4)} \mathcal{L}\{\sin(xt)\} dx$$

$$= \int_0^{\infty} \frac{1}{x(x^2+4)} \frac{x}{s^2+x^2} dx = \int_0^{\infty} \frac{1}{(x^2+4)(x^2+s^2)} dx$$

$$= \frac{1}{s^2-4} \int_0^{\infty} \left[ \frac{1}{x^2+4} - \frac{1}{x^2+s^2} \right] dx$$

$$= \frac{1}{s^2-4} \left[ \frac{1}{2} \tan^{-1} \frac{x}{2} - \frac{1}{s} \tan^{-1} \frac{x}{s} \right]_0^{\infty}$$

$$= \frac{1}{s^2-4} \left[ \frac{1}{2} \cdot \frac{\pi}{2} - \frac{1}{s} \cdot \frac{\pi}{2} \right] = \frac{\pi}{4(s^2-4)} \left[ 1 - \frac{2}{s} \right]$$

$$= \frac{\pi}{4(s^2-4)} \cdot \frac{s-2}{s} = \frac{\pi}{4s(s+2)} = \frac{\pi}{8} \left[ \frac{1}{s} - \frac{1}{s+2} \right]$$

$$f(t) = \frac{\pi}{8} [1 - e^{-2t}]$$

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Q.3 (15 points) Use Laplace transform to solve the heat equation

$$ku_{xx} = u_t, \quad 0 < x < L, \quad t > 0$$

under the following conditions

$$u(0, t) = 0, \quad u(L, t) = 0, \quad t > 0$$

and

$$u(x, 0) = x, \quad 0 < x < L$$

Apply residue theory and write final answer as a series.

Sol: Taking Laplace transform, we get

$$k \frac{d^2 U}{dx^2} = s U(x, s) - u(x, 0)$$

$$\frac{d^2 U}{dx^2} - \frac{s}{k} U = -\frac{x}{k}$$

$$u(0, t) = 0 \Rightarrow U(0, s) = 0 \quad \text{and} \quad u(L, t) = 0 \Rightarrow U(L, s) = 0$$

To find  $U_c$ :  $\frac{d^2 U}{dx^2} - \frac{s}{k} U = 0$ ,  $m = \pm \sqrt{\frac{s}{k}}$

$$U_c(x, s) = A \cosh \sqrt{\frac{s}{k}} x + B \sin \sqrt{\frac{s}{k}} x$$

To find  $U_p$ : Let  $U_p = A + Bx$ , then

$$-\frac{s}{k} (A + Bx) = -\frac{1}{k} x$$

$$\Rightarrow A = 0, \quad -\frac{s}{k} B = -\frac{1}{k} \Rightarrow B = \frac{1}{s}$$

$$U(x, s) = A \cosh \sqrt{\frac{s}{k}} x + B \sin \sqrt{\frac{s}{k}} x + \frac{x}{s}$$

$$U(0, s) = 0 \Rightarrow A = 0, \quad U(L, s) = 0 \Rightarrow B \sinh \sqrt{\frac{s}{k}} L + \frac{L}{s} = 0$$

$$B = -\frac{\frac{L}{s}}{\sinh \sqrt{\frac{s}{k}} L}$$

$$\therefore U(x, s) = -\frac{L}{s} \cdot \frac{\sinh \sqrt{\frac{s}{k}} x}{\sinh \sqrt{\frac{s}{k}} L} + \frac{x}{s}$$

$$s \sinh \sqrt{\frac{s}{k}} L$$

Check if  $s=0$  is a singularity or not?

$$U(x, s) = \frac{1}{s} \left[ x - \frac{L \left[ \sqrt{\frac{s}{k}} x + \frac{1}{3!} \left( \sqrt{\frac{s}{k}} x \right)^3 + \dots \right]}{\sqrt{\frac{s}{k}} L + \frac{1}{3!} \left( \sqrt{\frac{s}{k}} L \right)^3 + \dots} \right]$$

$$= \frac{1}{s} \left[ x - \frac{x + \frac{1}{3!} \frac{s}{k} x^3 + \dots}{1 + \frac{1}{3!} \frac{s}{k} L^2 + \dots} \right]$$

$$= \frac{1}{s} \left[ x - \left( x + \frac{s x^3}{6k} + \dots \right) \left( 1 + \frac{s L^2}{6k} + \dots \right) \right]$$

$$= \frac{1}{s} \left[ x - \left( x + \frac{s x^3}{6k} + \dots \right) \left( 1 - \frac{s L^2}{6k} - \dots \right) \right]$$

$$= \frac{1}{s} \left[ x - \left( x - \frac{s L^2 x}{6k} - \dots + \frac{s x^3}{6k} - \dots \right) \right]$$

$$= \frac{1}{s} \left[ \frac{s L^2 x}{6k} - \frac{s x^3}{6k} + O(s) \right]$$

$$= \frac{x L^2 - x^3}{6k} + O(s)$$

$s=0$  is a removable singularity.

$$\text{Now } \sinh \sqrt{\frac{s}{k}} L = 0 \Rightarrow \sqrt{\frac{s}{k}} L = n\pi i, \quad n=1, 2, \dots$$

$$s = -\frac{k n^2 \pi^2}{L^2}, \quad n=1, 2, 3, \dots \quad \text{Simple poles}$$

The residue of  $U(x, s) e^{st}$  at these poles

$$\lim_{s \rightarrow -\frac{k n^2 \pi^2}{L^2}} \left( s + \frac{k n^2 \pi^2}{L^2} \right) \frac{e^{st}}{s} \left( x - \frac{L \sinh \sqrt{\frac{s}{k}} x}{\sinh \sqrt{\frac{s}{k}} L} \right)$$

$$= \lim_{s \rightarrow -\frac{kn^2\pi^2}{L^2}} \frac{e^{st}}{s} \left[ \lim_{s \rightarrow -\frac{kn^2\pi^2}{L^2}} x \left( s + \frac{kn^2\pi^2}{L^2} \right) - \lim_{s \rightarrow -\frac{kn^2\pi^2}{L^2}} \frac{(s + \frac{kn^2\pi^2}{L^2}) L \operatorname{Sinh} \sqrt{\frac{s}{K}} x}{\operatorname{Sinh} \sqrt{\frac{s}{K}} L} \right]$$

$$= \frac{e^{-\frac{kn^2\pi^2}{L^2}t}}{-\frac{kn^2\pi^2}{L^2}} \left[ 0 - L \operatorname{Sinh} \frac{n\pi}{L} x i \cdot \lim_{s \rightarrow -\frac{kn^2\pi^2}{L^2}} \frac{s + \frac{kn^2\pi^2}{L^2}}{\operatorname{Sinh} \sqrt{\frac{s}{K}} L} \right]$$

$$= -\frac{L^3 e^{-\frac{kn^2\pi^2}{L^2}t}}{kn^2\pi^2} i \operatorname{Sin} \frac{n\pi}{L} x \cdot \lim_{s \rightarrow -\frac{kn^2\pi^2}{L^2}} \frac{1}{\frac{L}{2\sqrt{SK}} \operatorname{Cosh} \sqrt{\frac{s}{K}} L}$$

$$= \frac{-iL^3 e^{-\frac{kn^2\pi^2}{L^2}t}}{kn^2\pi^2} \operatorname{Sin} \frac{n\pi}{L} x \cdot \frac{1}{\frac{L}{2 \frac{kn\pi}{L} i} \operatorname{Cosh} \frac{n\pi}{L} L i}$$

$$= \frac{-iL^3 e^{-\frac{kn^2\pi^2}{L^2}t}}{kn^2\pi^2} \cdot \operatorname{Sin} \frac{n\pi}{L} x \cdot \frac{2 \frac{kn\pi}{L} i}{L} \cdot \frac{1}{\operatorname{Cos} n\pi}$$

$$= \frac{2L e^{-\frac{kn^2\pi^2}{L^2}t}}{n\pi} \operatorname{Sin} \frac{n\pi}{L} x (-1)^n$$

$$\therefore u(x, t) = \frac{2L}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^n}{n} e^{-\frac{kn^2\pi^2}{L^2}t} \operatorname{Sin} \frac{n\pi}{L} x$$

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Q.4 (10 points) Solve the integral equation using Fourier transform

$$\int_{-\infty}^{\infty} e^{-4t^2} f(x-t) dt = e^{-3x^2}.$$

Sol:  $e^{-4x^2} * f(x) = e^{-3x^2}$

⇒ Applying Fourier transform

$$\sqrt{\frac{\pi}{4}} e^{-\frac{\alpha^2}{16}} F(\alpha) = \sqrt{\frac{\pi}{3}} e^{-\frac{\alpha^2}{12}}$$

$$\Rightarrow F(\alpha) = \frac{2}{\sqrt{3}} e^{\frac{\alpha^2}{16} - \frac{\alpha^2}{12}} = \frac{2}{\sqrt{3}} e^{-\frac{\alpha^2}{4} \left(\frac{1}{12}\right)}$$

$$f(x) = \frac{1}{2\pi} \frac{2}{\sqrt{3}} \int_{-\infty}^{\infty} e^{-\frac{\alpha^2}{4} \left(\frac{1}{12}\right)} e^{i\alpha x} d\alpha$$

$$= \frac{1}{\pi\sqrt{3}} \int_{-\infty}^{\infty} e^{-\frac{1}{48} (\alpha^2 - 48\alpha x i)} d\alpha$$

$$= \frac{1}{\pi\sqrt{3}} \int_{-\infty}^{\infty} e^{-\frac{1}{48} [\alpha^2 - 2\alpha(24xi) + (24xi)^2 - (24xi)^2]} d\alpha$$

$$= \frac{1}{\pi\sqrt{3}} \int_{-\infty}^{\infty} e^{-\frac{1}{48} [\alpha - 24xi]^2 + \frac{1}{48} (24xi)^2} d\alpha$$

$$= \frac{1}{\pi\sqrt{3}} e^{-12x^2} \int_{-\infty}^{\infty} e^{-\frac{1}{48} u^2} du,$$

$$u = \alpha - 24xi$$

$$= \frac{1}{\pi\sqrt{3}} e^{-12x^2} \cdot \frac{\sqrt{\pi}}{\sqrt{\frac{1}{48}}} = \frac{1}{\sqrt{3}} 4\sqrt{3} e^{-12x^2} = 4e^{-12x^2}$$

Q.5 (10 points) Use Fourier transform to solve the heat equation

$$ku_{xx} = u_t, \quad -\infty < x < \infty, \quad t > 0$$

under the following conditions  $u(x, 0) = f(x)$  where  $f(x) = \begin{cases} u_0 & |x| \leq 1 \\ 0 & |x| > 1 \end{cases}$ .

Sol: Taking Fourier transform, we get

$$-k\alpha^2 U(\alpha, t) = \frac{dU}{dt} \Rightarrow U(\alpha, t) = A e^{-k\alpha^2 t}$$

$$u(x, 0) = f(x) \Rightarrow U(\alpha, 0) = \int_{-\infty}^{\infty} f(x) e^{-i\alpha x} dx$$

$$U(\alpha, 0) = \int_{-1}^1 u_0 e^{-i\alpha x} dx = \frac{u_0}{-i\alpha} e^{-i\alpha x} \Big|_{-1}^1$$

$$= -\frac{u_0}{i\alpha} (e^{-i\alpha} - e^{i\alpha}) = \frac{2u_0}{\alpha} \left[ \frac{e^{i\alpha} - e^{-i\alpha}}{2i} \right]$$

$$= \frac{2u_0}{\alpha} \sin \alpha$$

$$\Rightarrow A = \frac{2u_0}{\alpha} \sin \alpha$$

$$U(\alpha, t) = \frac{2u_0}{\alpha} \sin \alpha e^{-k\alpha^2 t}$$

$$u(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} 2u_0 \frac{\sin \alpha}{\alpha} e^{-k\alpha^2 t} e^{i\alpha x} d\alpha$$

$$= \frac{u_0}{\pi} \int_{-\infty}^{\infty} \frac{\sin \alpha}{\alpha} e^{-k\alpha^2 t} [\cos \alpha x + i \sin \alpha x] d\alpha$$

$$= \frac{u_0}{\pi} \int_{-\infty}^{\infty} \frac{\sin \alpha}{\alpha} \cos \alpha e^{-k\alpha^2 t} d\alpha + 0$$

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Q.6 (10 points) Solve using Mellin transform

$$x^2 u_{xx} + x u_x + y_{yy} = 0, \quad 0 \leq x < \infty, \quad 0 < y < 1$$

$$\text{under the conditions } u(x, 0) = 0 \text{ and } u(x, 1) = \begin{cases} A & 0 < x < 1 \\ 0 & x > 1 \end{cases}$$

Sol.: Taking Mellin transform, we get

$$p(1+p) \tilde{u}(p, y) - p \tilde{u}(p, y) + \frac{d^2 \tilde{u}}{dy^2} = 0$$
$$\Rightarrow \frac{d^2 \tilde{u}}{dy^2} + p^2 \tilde{u} = 0$$

$$\tilde{u}(p, y) = C_1 \cos py + C_2 \sin py$$

$$u(x, 0) = 0 \Rightarrow \tilde{u}(p, 0) = 0 \Rightarrow C_1 = 0$$

$$u(x, 1) = A H(1-x) \Rightarrow \tilde{u}(p, 1) = \int_0^{\infty} x^{p-1} A H(1-x) dx$$

$$\tilde{u}(p, 1) = A \int_0^1 x^{p-1} dx = \frac{A}{p} x^p \Big|_0^1 = \frac{A}{p}$$

$$\Rightarrow C_2 \sin p = \frac{A}{p} \Rightarrow C_2 = \frac{A}{p \sin p}$$

$$\tilde{u}(p, y) = \frac{A}{p \sin p} \sin py$$

$$u(x, y) = \frac{A}{2\pi i} \int_{c-i\infty}^{c+i\infty} x^{-p} \frac{\sin py}{p \sin p} dp$$

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Q.7 (10 points) Show the Hankel transform

$$\mathcal{H}_0\{e^{-ar}\} = \frac{a}{(a^2 + \alpha^2)^{3/2}}$$

Sol.  $\mathcal{H}_0\{e^{-ar}\} = \int_0^{\infty} r J_0(\alpha r) e^{-ar} dr$

$$= \int_0^{\infty} \frac{t}{\alpha} J_0(t) e^{-\frac{a}{\alpha} t} \frac{1}{\alpha} dt, \quad \alpha r = t$$
$$= \frac{1}{\alpha^2} \int_0^{\infty} t J_0(t) e^{-st} dt, \quad s = \frac{a}{\alpha}$$
$$= \frac{1}{\alpha^2} \mathcal{L}\{t J_0(t)\}$$

Since  $\mathcal{L}\{t^{\nu} J_{\nu}(t)\} = \frac{2^{\nu} \Gamma(\nu + \frac{1}{2})}{\sqrt{\pi} (1+s^2)^{\nu + \frac{1}{2}}}$

for  $\nu=0$ ,  $\mathcal{L}\{J_0(t)\} = \frac{\Gamma(\frac{1}{2})}{\sqrt{\pi} (1+s^2)^{\frac{1}{2}}} = \frac{1}{(1+s^2)^{\frac{1}{2}}}$ ,  $\Gamma(\frac{1}{2}) = \sqrt{\pi}$

$$\mathcal{H}_0\{e^{-ar}\} = \frac{1}{\alpha^2} (-1) \frac{d}{ds} \left\{ \mathcal{L}\{J_0(t)\} \right\}$$

$$= -\frac{1}{\alpha^2} \frac{d}{ds} \left( \frac{1}{(1+s^2)^{\frac{1}{2}}} \right) = -\frac{1}{\alpha^2} \left(-\frac{1}{2}\right) (1+s^2)^{-\frac{3}{2}} \cdot 2s$$

$$= \frac{1}{\alpha^2} \frac{s}{(1+s^2)^{3/2}} = \frac{1}{\alpha^2} \frac{\frac{a}{\alpha}}{\left(1 + \frac{a^2}{\alpha^2}\right)^{3/2}} = \frac{a}{(a^2 + \alpha^2)^{3/2}}$$

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Q.8 (10 points) Solve the axisymmetric biharmonic equation using Hankel transform

$$\begin{aligned}\nabla^4 u(r, z) &= 0, \quad 0 \leq r < \infty, \quad z > 0 \\ u(r, 0) &= f(r), \quad 0 \leq r < \infty, \\ \frac{\partial u}{\partial z} &= 0 \text{ on } z=0, \quad 0 \leq r < \infty \\ u(r, z) &\rightarrow 0 \text{ as } z \rightarrow \infty.\end{aligned}$$

Sol.  $\nabla^2 u = \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{\partial^2 u}{\partial z^2}$

$$\begin{aligned}\nabla^4 u &= \nabla^2 \nabla^2 u = \left( \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{\partial^2}{\partial z^2} \right) \left( \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{\partial^2}{\partial z^2} \right) u \\ &= \left( \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{\partial^2}{\partial z^2} \right)^2 u\end{aligned}$$

$$\mathcal{H}_0 \{ u(r, z) \} = \tilde{u}_0(\alpha, z)$$

$$\text{and } \mathcal{H}_0 \left\{ \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} \right\} = -\alpha^2 \tilde{u}_0(\alpha, z)$$

$$\Rightarrow \left( \frac{d^2}{dz^2} - \alpha^2 \right) \tilde{u}_0(\alpha, z) = 0, \quad z > 0$$

Roots of the auxiliary equation are  $\pm\alpha, \pm\alpha$

$$\tilde{u}_0(\alpha, z) = (A + Bz)e^{-\alpha z} + (C + Dz)e^{\alpha z}$$

The solution is bounded  $\Rightarrow$  Put  $C = D = 0$

$$\text{So } \tilde{u}_0(\alpha, z) = (A + Bz)e^{-\alpha z}$$

$$\text{Now } u(r, 0) = f(r) \Rightarrow \tilde{u}_0(\alpha, 0) = \tilde{f}(\alpha)$$

$$\frac{\partial u}{\partial z}(r, 0) = 0 \Rightarrow \frac{d}{dz} \tilde{u}_0(\alpha, 0) = 0$$

$$\frac{d\tilde{u}_0}{dz} = -\alpha(A + Bz)e^{-\alpha z} + Be^{-\alpha z}$$

$$\tilde{u}_0(\alpha, 0) = \tilde{f}(\alpha) \Rightarrow A = \tilde{f}(\alpha)$$

$$\frac{d\tilde{u}_0(\alpha, 0)}{d\alpha} = 0 \Rightarrow -\alpha A + B = 0 \Rightarrow B = \alpha \tilde{f}(\alpha)$$

$$\tilde{u}_0(\alpha, z) = \tilde{f}(\alpha) (1 + \alpha z) e^{-\alpha z}$$

Taking inverse Hankel transform, we get

$$u(r, z) = \int_0^{\infty} \alpha J_0(\alpha r) \tilde{f}(\alpha) (1 + \alpha z) e^{-\alpha z} d\alpha$$

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Q.9 (15 points) Solve the integral equation using Wiener-Hopf technique

$$\int_0^{\infty} e^{-|x-\xi|} u(\xi) d\xi = -\frac{1}{4}u(x) + 1, \quad 0 < x < \infty.$$

$u(x)$  is bounded as  $x \rightarrow \infty$ .

Sol. Note that  $u(x) = 4 - 4 \int_0^{\infty} e^{-|x-\xi|} u(\xi) d\xi$  is bounded

as  $x \rightarrow \infty$ . Thus  $u_+(\alpha)$  is analytic function of  $\alpha$  for  $\text{Im}(\alpha) > 0$ .

Here  $K(x-\xi) = e^{-|x-\xi|}$ ,  $\mu = -\frac{1}{4}$ ,  $f(x) = 1$

$0 < x < \infty$

Define  $f_+(x) = \begin{cases} 1 & x > 0 \\ 0 & x < 0 \end{cases}$

$$\text{and } f_+^*(\alpha) = \int_{-\infty}^{\infty} f_+(x) e^{-i\alpha x} dx = \int_0^{\infty} e^{-i\alpha x} dx$$

$$= -\frac{1}{i\alpha} e^{-i\alpha x} \Big|_0^{\infty} = \frac{1}{i\alpha} = -\frac{i}{\alpha},$$

$$K(x) = e^{-|x|} \text{ and } \mathcal{F}\{K(x)\} = \frac{2}{1+\alpha^2} = K^*(\alpha)$$

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To extend the problem to  $(-\infty, \infty)$ , we write

$$\int_{-\infty}^{\infty} e^{-|x-\zeta|} u(\zeta) d\zeta = \begin{cases} -\frac{1}{4} u(x) + f(x) & 0 < x < \infty \\ g(x) & -\infty < x < 0 \end{cases}$$

where  $u(x) = 0$  and  $f(x) = 0$  for  $x < 0$

$$g(x) = 0 \text{ for } x > 0$$

So  $u(x) = u_+(x)$ ,  $f(x) = f_+(x)$ ,  $g(x) = g_-(x)$

$$\text{and } \int_{-\infty}^{\infty} e^{-|x-\zeta|} u_+(\zeta) d\zeta = -\frac{1}{4} u_+(x) + 1 + g_-(x)$$

$$\mathcal{F} \Rightarrow \frac{2}{1+\alpha^2} u_+^*(\alpha) = -\frac{1}{4} u_+^*(\alpha) - \frac{i}{\alpha} + g_-^*(\alpha)$$

$$\frac{9+\alpha^2}{4(1+\alpha^2)} u_+^*(\alpha) = -\frac{i}{\alpha} + g_-^*(\alpha)$$

$$\frac{(\alpha+3i)(\alpha-3i)}{4(\alpha+i)(\alpha-i)} u_+^*(\alpha) = -\frac{i}{\alpha} + g_-^*(\alpha)$$

$$\frac{\alpha+3i}{4(\alpha+i)} u_+^*(\alpha) = \frac{-i(\alpha-i)}{\alpha(\alpha-3i)} + \frac{\alpha-i}{\alpha-3i} g_-^*(\alpha)$$

$$\frac{-i(\alpha-i)}{\alpha(\alpha-3i)} = \frac{-i}{3\alpha} - \frac{2i}{3(\alpha-3i)} \quad \text{using partial fractions}$$

$P_+(\alpha) = \frac{-i}{3\alpha}$  is analytic for  $\text{Im}(\alpha) < 0$

$P_-(\alpha) = \frac{-2i}{3(\alpha-3i)}$  is analytic for  $\text{Im}(\alpha) < 3$

$$\frac{\alpha+3i}{4(\alpha+i)} u_+^*(\alpha) + \frac{i}{3\alpha} = \frac{-2i}{3(\alpha-3i)} + \frac{\alpha-i}{\alpha-3i} g_-^*(\alpha)$$

LHS is analytic in the upper half plane  $\text{Im}(\alpha) > 0$

RHS is analytic in the lower half plane  $\text{Im}(\alpha) < 3$

By the Liouville's theorem, both sides define an

entire function  $E^*(\alpha)$  and LHS  $\rightarrow 0$  as  $\alpha \rightarrow \alpha$

and RHS  $\rightarrow 0$  as  $\alpha \rightarrow -\infty$ .

Thus  $E^*(\alpha) \equiv 0$

$$\frac{\alpha+3i}{4(\alpha+i)} U_+^*(\alpha) + \frac{i}{3\alpha} = 0$$

$$U_+^*(\alpha) = \frac{-4i(\alpha+i)}{3\alpha(\alpha+3i)}$$

$$\begin{aligned} U(x) &= \frac{1}{2\pi} \int_{ai-\infty}^{ai+\infty} U_+^*(\alpha) e^{i\alpha x} d\alpha \\ &= -\frac{2}{3\pi} i \int_{ai-\infty}^{ai+\infty} \frac{\alpha+i}{\alpha(\alpha+3i)} e^{i\alpha x} d\alpha \end{aligned}$$

Poles of  $h(\alpha) = \frac{\alpha+i}{\alpha(\alpha+3i)}$  are at  $\alpha=0$ ,  $\alpha=-3i$

$$\text{Res}[h(\alpha), \alpha=0] = \lim_{\alpha \rightarrow 0} \frac{\alpha+i}{\alpha+3i} e^{i\alpha x} = \frac{1}{3}$$

$$\text{Res}[h(\alpha), \alpha=-3i] = \lim_{\alpha \rightarrow -3i} \frac{\alpha+i}{\alpha} e^{i\alpha x} = \frac{2}{3} e^{3x}$$

$$U(x) = -\frac{2i}{3\pi} \cdot 2\pi i \left[ \frac{1}{3} + \frac{2}{3} e^{3x} \right] = \frac{4}{9} (1 + 2e^{3x})$$

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